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**PENETRATION OF THE EARTH'S FREE
OSCILLATIONS AT 54 MINUTE PERIOD INTO
THE ATMOSPHERE**

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Abstract. It is known that the fundamental spheroidal mode ${}_0S_2$ of the Earth free oscillation with the period of about 54 min forces atmospheric oscillations. We present a certain phase relationship for components of the ${}_0S_2$ multiplet, which is based on synchronous collocated microbarograph and seismograph observations. This relationship is both the first observational manifestation of the Pekeris mode of global atmospheric oscillations with the 54 min period and a further proof of penetrating the Earth's ${}_0S_2$ mode into the atmosphere. We show that the linear non-dissipative model of steady forced oscillations in isothermal atmosphere at rest does not describe penetrating the ${}_0S_2$ mode into the atmosphere adequately.

1 Introduction

There are Earth free oscillations observed with seismographs (e.g. Aki and Richards, 1980). The lowest frequency mode of the oscillations is the fundamental spheroidal mode ${}_0S_2$ with period of about 54 min. The Earth's departure from spherical symmetry such as diurnal rotation, ellipticity, and lateral heterogeneity splits the mode ${}_0S_2$ into 5 components spaced about 50 s apart. The first observations of these component occurred after very large earthquakes with a magnitude $M > 7$ (e.g. Buland *et al.*, 1979; Aki and Richards, 1980). However, subsequently they were also observed on seismically quiet days (e.g. Lin'kov *et al.*, 1991), much like the Earth free oscillations ${}_0S_n$ with $n \approx 15-65$ of the 2-8 min period range (Suda *et al.*, 1998; Tanimoto *et al.*, 1998).

Oscillations with period of about 54 min and even the splitting inherent in the ${}_0S_2$ mode have been observed in geomagnetic disturbances (Roldugin, 1967; Bobova *et al.*, 1990). The amplitude of these oscillations is intensified during global seismic activations. The components of the ${}_0S_2$ mode have also been detected with microbarograph in ground pressure variations (Garmash *et al.*, 1989; Lin'kov *et al.*, 1991). There are microbarographic observations of penetrating the ${}_0S_2$ oscillations into the atmosphere at seismically inactive period.

This paper constitutes a further observational study of forcing atmospheric oscillations by the Earth's ${}_0S_2$ oscillations. We use synchronous collocated measurements by a microbarograph and seismograph to determine a phase

difference between the seismic oscillations and the pressure ones for the different components of the ${}_0S_2$ mode. By comparing these phase differences, we first experimentally reveal the oscillation frequency corresponding to the so-called Pekeris solution in the global oscillation theory for the Earth's atmosphere (Dikii, 1965). We also discuss mechanisms of penetrating the ${}_0S_2$ mode into the atmosphere.

2 Instrumentation

The ground pressure (p) variations have been measured with a microbarograph, made on the basis of the standard millibarograph M-22 (Kozhevnikova *et al.*, 1980). As far as the seismograph is concerned, Kirnos' vertical pendulum (Savarensky and Kirnos, 1955), protected from the direct effect of variations of the atmospheric pressure and temperature with a special chamber (Lin'kov *et al.*, 1982), has been used. The sensors of both instruments are provided with converters which contain photodiodes. A set of filters, employed at the output of the converters, has resulted in the detection of oscillations with period (τ) in the ~ 0.5 – 5 h range. For this τ range the seismograph records vertical acceleration due to a joint action of ground displacement and gravitational acceleration perturbation. The microbarograph sensitivity at the converter output was in the 1 – 5 mWPa $^{-1}$ range, which has enabled us to measure pressure variations as small as some tenths of a μ bar. Amplification of the seismometric channel at $\tau = 1$ h equals about 10.

3 Determination of phase difference

A spheroidal mode of the Earth free oscillation is symbolized by ${}_kS_n$ (e.g. Buland *et al.*, 1979; Aki and Richards, 1980). Here k is the number of nodes between the surface and center of the Earth for the radial wavefunction of the mode. $k = 0$ gives the fundamental mode. The integer n is the degree of the spherical surface functions which describe the horizontal structure of motions associated with the mode. Namely, the radial (vertical) displacement for a m component of the ${}_kS_n$ mode splitting is proportional to the complex spherical harmonic $Y_n^m(\vartheta, \varphi)$ where ϑ is the co-latitude and φ is the longitude, $-n \leq m \leq n$. So, the ${}_kS_n$ mode is split into $2n + 1$ components (lines).

The frequencies of the ${}_0S_2$ mode components are 0.3000, 0.3048, 0.3095, 0.3140, 0.3185 mHz for $m = -2, -1, 0, 1, 2$, respectively (Buland *et al.*, 1979). The analysis of seismograph records, as a rule, makes possible to detect only some of the five components (e.g. Lin'kov *et al.*, 1991). The wide scatter of relative powers of the components in concrete records results mainly from interference between the ${}_0S_2$ multiplet components due to the small frequency spacing (e.g. Aki and Richards, 1980). It is established for the ${}_0S_2$ oscillations forced by a well-identified earthquake that the observed distribution of the Earth ${}_0S_2$ components over power depends on latitude and longitude of both a measurement site and earthquake source as well as on other parameters of the earthquake (e.g. Mendiguren, 1973; Stein and Geller, 1977). As for detecting the ${}_0S_2$ components with microbarographic measurements, the gravity waves of at-

atmospheric sources can hinder detection of these components. Moreover, it may turn out that some of the components are not detected due to a strong frequency dependence of the efficiency of oscillation transmission into the atmosphere (see Section 5).

Our goal is to compare the phase differences ($\Delta\psi$) between seismic and atmospheric oscillations for the ${}_0S_2$ components disposed on different sides from the central component ($m = 0$) along the frequency axis. We do succeed in finding three synchronous collocated microbarographic and seismographic runs to compare $\Delta\psi_m$ for $m > 0$ and $m < 0$ (Table 1). The sampling time of measurements is 4 min. As an example we show time series (Fig.1) and their power and coherence spectra (Figs 2 and 3) for one of the three cases. It can be seen (Fig.2) that the power spectra of seismic oscillations and ground pressure variations are significantly distinguished. However, the spectra exhibit common harmonics near the ${}_0S_2$ multiplet frequencies. Fig.3 demonstrates that the correlation between these harmonics of seismic and pressure spectra is statistically justified.

The achievement of the goal in hand is complicated by the facts that (i) the ${}_0S_2$ oscillations can be strongly unsteady and (ii) there is a need to separate the ${}_0S_2$ multiplet components closely spaced in frequency. The conventional fast Fourier transform is not suited for solving the problem. That is why we have used the method of record processing, described by Savina *et al.* (1985). First of all the raw data have been transformed to time series $F(t_i)$, where t is

the time, employing a numerical filter with bandpass around the ${}_0S_2$ multiplet frequencies. Further, $F(t)$ has been approximated by

$$F(t) = \sum_j (A_j \cos \sigma_j t + B_j \sin \sigma_j t) + F^*(t) \quad (1)$$

on the whole time interval of measurements, where A_j , B_j , and σ_j are the parameters of a j -harmonic and $F^*(t)$ is the residual "noise" function. The parameters were fitted in a least square sense with using the following procedure. The power spectrum of $F(t)$ gives the initial values of the parameters and the initial number of harmonics used at the fitting. Namely, the harmonics having amplitudes in excess of 1/3 of the largest amplitude are taken. The parameters are sequentially determined in decreasing order of the amplitude of harmonic. Those of each harmonic are evaluated from time series derived by eliminating the contributions to $F(t)$ of harmonics with larger amplitudes than the amplitude of a harmonic considered. This special feature of the procedure makes it possible to detect harmonics unresolved by a conventional spectral analysis. That is particularly important for the closely spaced ${}_0S_2$ components which interfere and can be manifested themselves in a power spectrum as a single wide asymmetric peak. As many cycles of parameter determination are made as it needs for the cessation of a change of the parameter values. The fitting procedure has resulted in decreasing the dispersion of the time series $F^*(t_i)$ to no more than 6% of that for the time series $F(t_i)$.

We have utilized the above procedure of approximating the function $F(t)$ to obtain the frequencies σ_j of the ${}_0S_2$ multiplet components for a subsequent

search of their phases. To derive $\Delta\psi_m$ reliably, we have checked phase behaviour to avoid any perceptible jumps of phase during our measurements. Time intervals (common to microbarographic and seismographic records) of the duration between 10 h and 30 h with a high signal-to-noise ratio and without the phase jumps have been found for each detected ${}^0\text{S}_2$ component. [See procedure details in Savina *et al.* (1985).] The oscillation phases in these intervals have been analyzed. The comparison of $\Delta\psi_m$ has been led to a surprising conclusion:

$$|\Delta\psi_{m>0} - \Delta\psi_{m<0}| = \pi \quad (2)$$

taking into account an experimental error (Table 1).

4 Model of global atmospheric oscillations

The model of atmospheric oscillations, forced by seismic ones, should be considered for understanding the observational inference in Eq. (2). Following Garmash *et al.* (1989), we present a brief description of a simple model for steady global atmospheric oscillations of the periods considered. The atmosphere is considered as a layer of small thickness compared with the radius of the Earth (a). The background atmosphere is assumed to be isothermal and at rest. Thus, the basic density ($\bar{\rho}$) and pressure (\bar{p}) are proportional to $\exp(-z/H)$, where z is the height, and H is the scale height. The gravitational potential perturbation (Ω), induced by seismic oscillations, is assumed to be independent of z . The atmospheric oscillations are described by a set of the

following linear equations neglecting of both dissipative processes and diurnal rotation:

$$\begin{aligned}
\frac{\partial u}{\partial t} &= -\frac{1}{a} \frac{\partial}{\partial \vartheta} \left(\frac{p'}{\bar{\rho}} + \Omega \right), \\
\frac{\partial v}{\partial t} &= -\frac{1}{a \sin \vartheta} \frac{\partial}{\partial \varphi} \left(\frac{p'}{\bar{\rho}} + \Omega \right), \\
\frac{\partial w}{\partial t} &= -\frac{1}{\bar{\rho}} \frac{\partial p'}{\partial z} - g \frac{\rho'}{\bar{\rho}}, \\
\frac{\partial \rho'}{\partial t} - \frac{w \bar{\rho}}{H} &= -\bar{\rho} \chi, \\
\frac{\partial p'}{\partial t} - \frac{w \bar{\rho}}{H} &= \gamma g H \left(\frac{\partial \rho'}{\partial t} - \frac{w \bar{\rho}}{H} \right),
\end{aligned} \tag{3}$$

where u , v , and w are the respective ϑ , φ , and z components of the velocity \mathbf{V} ; ρ' and p' are the density and pressure perturbations, respectively; g is the acceleration of gravity; γ is the ratio of the specific heat capacity at constant pressure to that at constant volume; $\chi \equiv \text{div} \mathbf{V}$. (The effect of the Earth's daily rotation is briefly considered in Section 5.)

For a steady periodic process any dependent variable (f) is written

$$f(t, \vartheta, \varphi, z) = f_\sigma(\vartheta, \varphi, z) e^{i\sigma t} \tag{4}$$

with σ the oscillation frequency. Using Eq. (4) in Eq. (3) we have

$$p'_\sigma = \frac{ig\bar{\rho}}{\sigma} (\gamma H \chi_\sigma - w_\sigma) \tag{5}$$

and

$$\left(\frac{\sigma^2}{g} + \frac{g}{\sigma^2} \Delta_{\vartheta, \varphi} \right) w_\sigma = \gamma H \left(\frac{1}{H} - \frac{\partial}{\partial z} + \frac{g}{\sigma^2} \Delta_{\vartheta, \varphi} \right) \chi_\sigma - \frac{i}{\sigma} \Delta_{\vartheta, \varphi} \Omega_\sigma, \tag{6}$$

where $\Delta_{\vartheta, \varphi}$ is the horizontal Laplace operator. The equations (3) are reduced to a single equation for χ_σ alone:

$$\left(H \frac{\partial^2}{\partial z^2} - \frac{\partial}{\partial z} + \frac{\sigma^2}{\gamma g}\right) \chi_\sigma = \frac{g}{\sigma^2} \Delta_{\vartheta, \varphi} \left[\left(\frac{\gamma - 1}{\gamma} - \frac{\sigma^2 H}{g} \right) \chi_\sigma + \frac{i\sigma^3}{\gamma g^2} \Omega_\sigma \right]. \quad (7)$$

Since a spherical surface function $Y_n^m(\vartheta, \varphi)$ is the solution of the equation

$$a^2 \Delta_{\vartheta, \varphi} Y_n^m + n(n+1) Y_n^m = 0, \quad (8)$$

any f_σ is written

$$f_\sigma(z, \vartheta, \varphi) = \sum_{n,m} f_{\sigma,n}^m(z) Y_n^m(\vartheta, \varphi). \quad (9)$$

Following the common practice of changing variables,

$$x = \frac{z}{H} \quad (10)$$

and

$$y_{\sigma,n}^m = \chi_{\sigma,n}^m e^{-\frac{x}{2}}, \quad (11)$$

the problem of global atmospheric oscillations is reduced to solving the vertical structure equation:

$$\frac{d^2 y_{\sigma,n}^m}{dx^2} - \lambda_{\sigma,n}^2 y_{\sigma,n}^m = \beta_{\sigma,n}^m e^{-\frac{x}{2}}, \quad (12)$$

where

$$\lambda_{\sigma,n}^2 = \frac{1}{4} - \frac{(\gamma - 1)H}{\gamma h_{\sigma,n}} + C_\sigma \left(\frac{H}{h_{\sigma,n}} - \frac{1}{\gamma} \right), \quad (13)$$

$$\beta_{\sigma,n}^m = -\frac{i\sigma^3 H}{\gamma g^2 h_{\sigma,n}} \Omega_{\sigma,n}^m, \quad (14)$$

$$h_{\sigma,n} = \frac{a^2 \sigma^2}{n(n+1)g}, \quad (15)$$

$$C_\sigma = \frac{H\sigma^2}{g}. \quad (16)$$

The upper boundary condition Eq. (12) is traditional, requiring the finiteness of $y_{\sigma,n}^m(x)$ as $x \rightarrow \infty$. The lower boundary condition at $x = 0$, like Ω , incorporates forcing atmospheric oscillations by seismic oscillations through the piston mechanism:

$$w_{\sigma,n}^m(0) = i\sigma \Delta a_{\sigma,n}^m, \quad (17)$$

where $\Delta a_{\sigma,n}^m$ is the amplitude of ground displacement for the oscillation mode (σ, nm) , and $w_{\sigma,n}^m$ may be written from Eq. (6) as

$$w_{\sigma,n}^m = \frac{\gamma}{\frac{1}{h_{\sigma,n}} - \frac{\sigma^2}{g}} \left[e^{\frac{x}{2}} \left(\frac{d}{dx} + \frac{H}{h_{\sigma,n}} - \frac{1}{2} \right) y_{\sigma,n}^m + \frac{\beta_{\sigma,n}^m}{C_\sigma} \right]. \quad (18)$$

Finally, the solution of the Eq. (12) is

$$y_{\sigma,n}^m = \frac{\beta_{\sigma,n}^m e^{-\frac{x}{2}}}{\frac{1}{4} - \lambda_{\sigma,n}^2} + \frac{e^{-\lambda_{\sigma,n}x}}{\frac{H}{h_{\sigma,n}} - \lambda_{\sigma,n} - \frac{1}{2}} \times \left\{ \frac{i\sigma}{\gamma} \left(\frac{1}{h_{\sigma,n}} - \frac{\sigma^2}{g} \right) \Delta a_{\sigma,n}^m + \beta_{\sigma,n}^m \left[\frac{1 - \frac{H}{h_{\sigma,n}}}{\frac{1}{4} - \lambda_{\sigma,n}^2} - \frac{1}{C_\sigma} \right] \right\}. \quad (19)$$

Using Eq. (5), we also give the expression for the pressure perturbation, which will be used in the discussion:

$$\begin{aligned} \frac{p'_{\sigma,n}{}^m}{\bar{p}} &= \frac{\Delta a_{\sigma,n}^m}{H} \left[1 - \frac{H}{\frac{H}{h_{\sigma,n}} - \lambda_{\sigma,n} - \frac{1}{2}} \left(\frac{1}{h_{\sigma,n}} - \frac{\sigma^2}{g} \right) \right] e^{(\frac{1}{2} - \lambda_{\sigma,n})x} + \\ &+ \frac{C_\sigma \Omega_{\sigma,n}^m}{gh_{\sigma,n}} \left\{ \frac{1}{\frac{1}{4} - \lambda_{\sigma,n}^2} + \frac{e^{(\frac{1}{2} - \lambda_{\sigma,n})x}}{\frac{H}{h_{\sigma,n}} - \lambda_{\sigma,n} - \frac{1}{2}} \left[\frac{1 - \frac{H}{h_{\sigma,n}}}{\frac{1}{4} - \lambda_{\sigma,n}^2} - \frac{1}{C_\sigma} \right] - \right. \\ &\left. - \frac{1}{H} \left[e^{(\frac{1}{2} - \lambda_{\sigma,n})x} - 1 \right] \left[\frac{1 - \frac{H}{h_{\sigma,n}}}{\frac{1}{4} - \lambda_{\sigma,n}^2} - \frac{1}{C_\sigma} \right] \right\}. \quad (20) \end{aligned}$$

The first and second terms of Eq. (20) describe piston and gravitational forcings of atmospheric oscillations, respectively.

5 Discussion

The detection of periodicity at 54 min in auroral bursts by Roldugin (1967) may be considered as the first evidence of penetrating the ${}_0S_2$ mode into the atmosphere. Later Bobova *et al.* (1990) have revealed all five components of the ${}_0S_2$ multiplet in the spectrum of the 1-min geomagnetic AE-index for the 14-year period. Both results may be interpreted in terms of the ionospheric wind dynamo (e.g. Kelley, 1989). That is the wind 54-min variation in the lower thermosphere generates a corresponding current variation which leads to resultant periodicity at 54 min in electric fields and magnetic perturbations. Other observational evidence of penetrating the Earth's ${}_0S_2$ mode into the atmosphere consists of revealing steady prominent features near the ${}_0S_2$ multiplet frequencies in spectra of ground pressure variations (Garmash *et al.*, 1989; Lin'kov *et al.*, 1991). The relation between seismic and atmospheric oscillations is confirmed, for example, by deriving the coherence coefficient of 0.72 (at the 95% confidence level) at the 55 min period from synchronous microbarographic and seismographic runs in St.Petersburg (60°N, 30°E) in 1981-82 (Garmash *et al.*, 1989).

Garmash *et al.* (1989) have put forward the piston mechanism for the explanation of penetrating long-period Earth's oscillations into the atmosphere. The penetrating can occur at the resonant frequencies determined by the condition

$$h_{\sigma,n} = \gamma H \tag{21}$$

for free oscillations of the atmosphere (Dikii, 1965). This condition provides

$$\frac{H}{h_{\sigma,n}} - \lambda_{\sigma,n} - \frac{1}{2} = 0. \quad (22)$$

Conceivably the piston mechanism of the condition (21) is realized for penetrating some Earth's oscillations into the atmosphere (Petrova *et al.*, 1996). This mechanism is yet unlikely suited for the generation of atmospheric oscillations by the Earth's ${}_0S_2$ mode. The point is that the horizontal structure of the ${}_0S_2$ mode corresponds to the $Y_n^m(\vartheta, \varphi)$ functions with $n = 2$, whereas near the 54 min period the condition (21) provides penetration of oscillations with $n = 35-40$.

However, there is another solution in the free oscillation theory for the atmosphere (Dikii, 1965):

$$h_{\sigma,n} = \frac{g}{\sigma^2}. \quad (23)$$

Following Dikii, we name this solution as the Pekeris solution. Substitution of Eq. (15) into Eq. (23) determines Pekeris' frequencies:

$$\sigma_{P,n} = \sqrt[4]{n(n+1)} \sqrt{\frac{g}{a}}. \quad (24)$$

The expression (20) shows the resonance amplification of $p'_{\sigma,n}$ for $x > 0$ near the frequencies $\sigma_{P,n}$ due to gravitational potential perturbation. We should mention in passing that Dikii (1965) considered free atmospheric oscillations at $\sigma_{P,n}$ as physically unrealizable ones.

Garmash *et al.* (1989) called attention to the following surprising coincidences for the second Pekeris mode ($n = 2$) and the Earth's ${}_0S_2$ mode. (i) The period

$\tau_{P,2} \approx 54$ min is practically equal to τ of the central component ($m = 0$) of the ${}_0S_2$ multiplet. (ii) Their horizontal structures agree as described by $Y_2^m(\vartheta, \varphi)$. It is these coincidences that points a way for the explanation of penetrating the ${}_0S_2$ mode into the atmosphere. Unfortunately, the simple oscillation model considered is not able to provide any quantitative treatment of the penetrating. Below we will revert to this question in detail.

The coincidental points specified also suggest explanation of the phase relationship Eq. (2). The explanation is based on the following:

(a) The gravitational potential perturbation and ground displacement result from common seismic processes. Therefore,

$$\Omega_{\sigma,n}^m = c \Delta a_{\sigma,n}^m \quad (25)$$

with c the unknown complex number. A simple example of the c expression was derived by Garmash et al. (1989) on the assumption that Ω is induced by small vertical displacements of the earth's surface neglecting earth's density variations:

$$c = \frac{3g}{2n+1} \left(1 + \frac{z}{a}\right)^{-(n+1)} \quad (26)$$

with g the acceleration of gravity at the surface.

(b) Relative contributions of ground displacement and gravitational acceleration perturbation to the acceleration, recorded by a seismograph, are unknown. However, from Eq. (25), the recorded acceleration is linear in the ground displacement with some certain complex factor.

(c) The expression (20) for the pressure perturbation contains terms with both

the factor of $(1/h_{\sigma,n} - \sigma^2/g)$ and its inverse.

(d) The expression $(1/h_{\sigma,n} - \sigma^2/g)$ changes its sign at passage of σ through $\sigma_{P,n}$. In the case under consideration that is easily seen from the expression

$$\frac{1}{h_{\sigma,2}} - \frac{\sigma^2}{g} \simeq -\frac{4\sqrt{6}}{a} \frac{\delta\sigma}{\sigma_{P,2}} \quad (27)$$

obtained for $|\delta\sigma| \ll \sigma_{P,2}$ where $\delta\sigma = \sigma - \sigma_{P,2}$.

Thus, one might therefore expect that $p'_{\sigma,2}{}^m$ changes sign at going from $\sigma_{m<0}$ to $\sigma_{m>0}$. This sign change must be revealed in observations as the phase relationship Eq. (2). The relationship may be interpreted as an observational proof of global atmospheric oscillations associated with the second Pekeris mode. It should be noted that the relationship (2) can be explained only under the assumption of forcing atmospheric oscillations by the Earth's ${}_0S_2$ oscillations. The fulfilment of the relationship is therefore a further proof of penetrating the ${}_0S_2$ mode into the atmosphere.

It should be stressed that the simple model of Section 4 gives only a general idea of the explanation of the relationship (2). The relationship can not be derived with this model, as the model is not suited to estimate properly the efficiency of forcing atmospheric oscillations by the Earth's ${}_0S_2$ oscillations. That is easily seen if $p'_{\sigma,2}{}^m$ is considered in the vicinity of $\sigma_{P,2}$ for $|\delta\sigma|/\sigma_{P,2}$ of no more than the value of $6 \cdot 10^{-2}$, corresponding the frequency interval of the ${}_0S_2$ multiplet (between the components with $m = -2$ and 2). Taking into account also that $C_\sigma \ll 1$ for $\sigma = \sigma_{P,2}$ ($C_\sigma \approx 3 \cdot 10^{-3}$), the formula (20) may

be approximated by the expression

$$\frac{p'_{\sigma,2}{}^m}{\bar{p}} \simeq \frac{\Delta a_{\sigma,2}^m}{H} \left[1 - \frac{4\sqrt{6}H}{a} \frac{\delta\sigma_m}{\sigma_{P,2}} \right] + \frac{\sqrt{6}\Omega_{\sigma,2}^m}{ag} \left\{ 1 + x \left[\frac{\sqrt{6}H}{2a} \left(\frac{\delta\sigma_m}{\sigma_{P,2}} \right)^{-1} + \frac{1}{\gamma} - 1 \right] \right\}. \quad (28)$$

We use also

$$\eta = \frac{|p'_{\sigma,n}{}^m|}{\Delta a_{\sigma,n}^m} \quad (29)$$

as an efficiency coefficient for forcing atmospheric oscillations by seismic oscillations (Garmash et al., 1989).

The "piston" term of Eq. (28) does not change its sign at passage of σ through $\sigma_{P,2}$, since

$$\frac{4\sqrt{6}H}{a} \frac{|\delta\sigma|}{\sigma_{P,2}} \ll 1. \quad (30)$$

In the case of the piston mechanism for the ${}_0S_2$ mode

$$\eta = \frac{\bar{p}}{H}. \quad (31)$$

This small η does not make possible to detect an effect of the mechanism by microbarographs (Garmash *et al.*, 1989).

As to the gravitational term of Eq. (28), the resonance turns out very "sharp" to the extent that it can provide penetrating into the atmosphere only for the central component ($m = 0$) of the ${}_0S_2$ multiplet. Using, for example, c from (26) for $z = 0$ and $n = 2$,

$$\eta = \frac{3\sqrt{6}\bar{p}}{5a} \left| 1 - x \left(1 - \frac{1}{\gamma} \right) \right| \quad (32)$$

outside the resonance peak. Beyond this peak, the efficiency of gravitational transmission of the ${}_0S_2$ components into the atmosphere is much less than that

by the piston mechanism because of $H \ll a$. There is little likelihood that taking account of the earth's density variations in Ω can enhance the gravitational η by 2–3 orders which are required for equalizing the contributions of both mechanisms to η .

Now consider whether the inferences derived here are changed by introducing the Earth's daily rotation to the model of global atmospheric oscillations. The rotation splits frequencies of free atmospheric oscillations (e.g. Dikii, 1965). We have estimated the order-of-magnitude of the splitting basing on the simple linear model of global oscillations in the limit of large $\sigma/2\omega$ where ω is the angular velocity of Earth [see Dikii's formula (2.8)]. For free atmospheric oscillations described with using the functions $Y_n^m(\vartheta, \varphi)$ the rotational perturbation of the frequency is

$$\Delta\sigma_n^m = \frac{m\omega}{n(n+1)}, \quad (33)$$

$-n \leq m \leq n$. For the second atmospheric Pekeris mode the frequency spacing between neighboring multiplet components is $\omega/6 = 1.93 \mu\text{Hz}$. That for the ${}_0\text{S}_2$ multiplet is 4.5-4.8 μHz . Taking into account that a m-component of the ${}_0\text{S}_2$ mode penetrates into the atmosphere by forcing atmospheric oscillations in the like m-component of the Pekeris mode two conclusions can be made. (i) The spacing between the ${}_0\text{S}_2$ components is more than double that between the Pekeris components. At going from $\sigma_{m<0}$ to $\sigma_{m>0}$ the sign change in $p_{\sigma,2}^m$ is therefore provided, and the explanation of the phase relationship Eq. (2) retains correspondingly its validity. (ii) Due to the Earth's rotation the ${}_0\text{S}_2$ multiplet frequencies turn out closer to the respective resonant frequencies of

the Pekeris mode. This fact is capable of increasing the efficiency of gravitational transmission of the ${}_0S_2$ components into the atmosphere.

So, the simple model considered does not provide the ultimate answer what mechanism forces atmospheric oscillations at the frequencies of the ${}_0S_2$ multiplet. There is some preference of the gravitational mechanism, owing to the "sharp" resonance. However, it may turn out in reality that the piston mechanism contributes as well. Further progression the problem may be made by introducing the following improvements into the model:

- (a) The expression for Ω should take into account the earth's density variations.
- (b) The Earth's rotation, dissipative processes, and real fields of temperature and wind should be incorporated into the atmospheric oscillation model.
- (c) The main what should be done is to turn to the problem of unsteady oscillations. As a source of forcing, the Earth's ${}_0S_2$ oscillations are variable, being completely damped out from time to time. Therefore, the steady oscillation of the simple model is unlikely to be realized.

6 Conclusion

The synchronous collocated microbarograph and seismograph measurements have resulted in the phase relationship Eq. (2) for the ${}_0S_2$ multiplet components with different signs of the number m . For the first time the second Pekeris mode of global atmospheric oscillations has been revealed experimen-

tally. The relationship should also be read as a further proof of penetrating the Earth's 0S_2 mode into the atmosphere. The linear non-dissipative model of steady forced oscillations in isothermal atmosphere at rest does not describe penetrating the 0S_2 mode into the atmosphere adequately.

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FIGURE CAPTIONS

Fig.1. Synchronous microbarographic (B) and seismographic (S) records during 31 January – 3 February 1982 at St.Petersburg (60°N , 30°E), in relative units. The upper time series are raw data. The lower time series have been derived from the upper ones by eliminating long-period trends using a smoothing filter.

Fig.2. Power spectra of the unfiltered microbarographic (B) and seismographic (S) time series shown in Fig.1. Two prominent peaks in both spectra, attributed to the ${}_{0}\text{S}_2$ multiplet components, are marked off by a bar.

Fig.3. Coherence spectrum of the filtered microbarographic and seismographic time series shown in Fig.1. The hypothesis of statistical independence of seismic and ground pressure oscillations from each other is rejected with 95% and 99% probability for coherence values above the solid and dashed horizontal lines, respectively. The left maximum corresponds to the 57.1 min period.